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Metamorphic $In_{0.7}Al_{0.3}As/In_{0.69}Ga_{0.31}As$ thermophotovoltaic devices grown on graded $InAs_yP_{1-y}$ buffers by molecular beam epitaxy

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ABSTRACT

The application of an $In_{0.7}AI_{0.3}As$ window layer material on the properties of lattice-mismatched (LMM) single-junction (SJ) $In_{0.69}Ga_{0.31}As$ thermophotovoltaic (TPV) cells grown by solid source molecular beam epitaxy (MBE) was studied as an alternative to conventional InAsP window layers. For these 0.60 eV bandgap SJ devices, high performance was achieved, displaying open-circuit voltage of 355 mV, power density of 0.532 W/cm² and a fill factor of 66.5% measured at a current density value of 2.26 A/cm². The measured internal quantum efficiencies were close to ~100% at wavelengths between approximately 1.2–1.4 mm, suggestive that low carrier recombination rates in the vicinity of this alternative $In_{0.70}$ $AI_{0.30}As/In_{0.69}Ga_{0.31}As$ window/emitter interface has been achieved. The quantum efficiency data coupled with the device parameters show that InAlAs can be successfully substituted for conventional InAsP window/emitter material pairing as opposed to the conventional InAsP/InGaAs window/emitter design, this approach avoids the problematic III-P/III-As interface transition for MBE growth of TPV devices that has been shown to impact carrier collection efficiency [Hudait MK, Lin Y, Palmisiano MN, Ringel SA. 0.6-eV band gap $In_{0.69}Ga_{0.31}As$ thermophotovoltaic devices grown on $InAs_P/I_{1-y}$ step-graded buffers by molecular beam epitaxy. IEEE Electron Dev Lett 2003;24(9):538–40].

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1. Introduction

In_xGa_{1-x}As based thermophotovoltaic (TPV) devices grown on InP substrates are of interest for a variety of terrestrial and space energy conversion applications [1–7]. For high conversion efficiency and power density, InGaAs TPV cells with band gaps from 0.50 eV to 0.74 eV are required. To achieve these band gaps requires an In content (x) of the active $In_xGa_{1-x}As$ TPV layers well in excess of the 53% composition that provides a convenient lattice match to InP substrates. The subsequent lattice mismatch between the $In_xGa_{1-x}As$ TPV device and the InP substrate (1.1% for 0.6 eV bandgap TPV cells) necessitates a buffer scheme to reduce the threading dislocation density that would otherwise propagate through the device layers. Also, the emitter bulk material and interface quality between the internally lattice-matched emitter and window layers are of great importance for achieving high-performance TPV cells, as these regions impact minority carrier collection efficiency [8,9], in particular at shorter wavelengths.

TPV cells based on $In_xGa_{1-x}As$ can utilize one of several window layer materials. InAsP is commonly adopted to provide a classic, symmetrical InAsP/InGaAs/InAsP double heterostructure design since InAsP is also used for both the back surface field and graded buffer layers in common metal-organic chemical vapor deposition (MOCVD) grown TPV cells [5-7,10]. However, solid source molecular beam epitaxy (MBE) is now receiving interest for TPV applications due to its extreme precision and growth uniformity [2], and since it provides an opportunity to investigate and potentially optimize lattice-mismatched (LMM) TPV structures within a very different growth regime compared to MOCVD. On the other hand, for TPV structures grown using MBE, the use of InAsP windows on InGaAs TPV devices is complicated by the need for a short growth interruption at the InAsP/InGaAs window/emitter interface. This is required to establish optimum arsenic and phosphorus beam fluxes prior to growth of the thin $InAs_yP_{1-y}$ window layer that must be lattice matched to the $In_xGa_{1-x}As$ emitter. We have demonstrated the sensitivity of local carrier recombination to the interface switching conditions for InP/InGaAs interfaces [11]. Typical 0.6 eV bandgap InGaAs TPV cells require an InAs_{0.32}P_{0.68} composition window layer to maintain lattice matching, resulting in a nominal bandgap of $\sim 1 \text{ eV}$ for the passive window layer [12,13]. In contrast, the $In_xAl_{1-x}As$ material system, which has been relatively unexplored for TPV applications, also provides latticematched window choices for $In_xGa_{1-x}As$ TPV devices, and for the





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Fig. 1. Nomarski image of an $In_{0.69}Ga_{0.31}As$ TPV device structure grown on InP substrate using $InAs_yP_{1-y}$ step-graded buffer.

same 0.6 eV InGaAs bandgap cell, a lattice-matched $In_{0.7}AI_{0.3}As$ would also provide a nominal window layer bandgap of ~1 eV [12]. From the perspective of MBE growth, an InAlAs/InGaAs window/emitter structure within a TPV cell would avoid the need for a growth interruption at the crucial window/emitter interface, which may improve the short wavelength performance reported for earlier MBE-grown TPV devices [2].

In this study, the effect of using an $In_{0.7}AI_{0.3}As$ window layer on the properties of lattice-mismatched (LMM) single-junction (SJ) n/p $In_{0.69}Ga_{0.31}As$ TPV cell has been investigated. The LMM $In_{0.69}Ga_{0.31}As$ TPV device structures grown by MBE were analyzed by cross-sectional transmission electron microscopy (XTEM), triple crystal X-ray diffraction and Nomarski microscopy to evaluate defect properties, strain relaxation and surface morphology, respectively. Single junction LMM In_{0.69}Ga_{0.31}As TPV devices with band gaps of 0.6 eV were processed and high performance is achieved for these cells implies In_{0.7}Al_{0.3}As layer can be used as an alternative window layer material for In_{0.69}Ga_{0.31}As TPV device structures.

2. MBE growth and device processing

Lattice-mismatched SJ In_{0.69}Ga_{0.31}As TPV structures were grown on semi-insulating 2° offcut (100) InP substrates using a solid source MBE system equipped with valved cracker sources for arsenic and phosphorus. InP substrate oxide desorption was performed at 510 °C under a phosphorus overpressure of $\sim 1 \times 10^{-5}$ Torr and verified by observing a strong (2 \times 4) reflection high-energy electron diffraction (RHEED) pattern, indicating a clean (100) InP surface. An undoped 0.2 µm InP buffer layer was then deposited under a stabilized P₂ flux prior to the growth of an n-type Si-doped $InAs_{\nu}P_{1-\nu}$ step-graded buffer. The step-graded $InAs_{\nu}P_{1-\nu}$ buffer consisted of 4-steps, with the final compositions of InAs_{0.32}P_{0.68} providing a lattice-matched "virtual" substrate for In_{0.69}Ga_{0.31}As TPV overgrowth. The surface morphology of LMM $In_{0.69}Ga_{0.31}As$ TPV device grown on InP substrate using $InAs_{\nu}P_{1-\nu}$ step-graded buffer shown in Fig. 1 exhibits a well-developed 2D crosshatch pattern and this uniform crosshatch pattern is an indication of metamorphic device layer growth. Reciprocal space maps (RSMs) using triple crystal X-ray diffraction were used to determine the alloy composition, strain relaxation and the internally lattice-matched condition to the upper layer of $InAs_{\nu}P_{1-\nu}$ step-graded buffer InAs_{0.32}P_{0.68}, In_{0.69}Ga_{0.31}As device layer and In_{0.7}Al_{0.3}As window layer. Fig. 2a and b shows the RSMs for (004) and (115)



Fig. 2. (a) Symmetric (004) and (b) asymmetric (115) reciprocal space maps of an $In_{0.7}Al_{0.3}As/In_{0.69}Ga_{0.31}As/InAs_{0.32}P_{0.68}$ TPV device structure grown on (001) InP substrate using $InAs_yP_{1-y}$ step-graded buffer obtained using an incident X-ray beam oriented along the [110] direction.

reflections from LMM $In_{0.69}Ga_{0.31}As$ TPV device structure grown on InP substrate with the incident beam is along the [110] direction. In Fig. 2a, the (004) reciprocal space maps exhibit five distinct reciprocal lattice units (RLP) which corresponds to the InP substrate, three steps in $InAs_yP_{1-y}$ graded buffer and the lattice-matched $In_{0.69}Ga_{0.31}As$ layer. The $In_{0.7}Al_{0.3}As/In_{0.69}Ga_{0.31}As$ layer RLP lies directly on the $InAs_{0.32}P_{0.68}$ buffer layer, thus in-plane

lattice constant of the In_{0.69}Ga_{0.31}As layer is coherent with the uppermost buffer layer, indicating a close internally latticematched condition. The degree of relaxation of the In_{0.69}Ga_{0.31}As overlayer is found to be more than 90% along with the measured values of in-plane lattice constant (from Fig. 2b) and out-of-plane lattice constant (from Fig. 2a). Full details on the growth and properties of the InAs_yP_{1-y} step-graded buffer and the minority carrier



b



Fig. 3. (a) Schematic cross-section and (b) cross-sectional TEM micrograph of a typical lattice-mismatched In_{0.69}Ga_{0.31}As n/p/n TPV test structure grown with an In_{0.7}Al_{0.3}As window layer and InAs_yP_{1-y} step-graded buffer on InP substrate.

а

1.2

1.0

0.8

0.6

lifetime of In_{0.69}Ga_{0.31}As on top of this InAs_vP_{1-v} step-graded buffer were previously reported [3,14,15].

The schematic cross-section of LMM n/p/n In_{0.69}Ga_{0.31}As TPV structure shown in Fig. 3a with an In_{0.7}Al_{0.3}As window layer allows the use of the desired n-on-p cell configuration with an n-type In-As_{0.32}P_{0.68} lateral conduction layer (LCL) to interconnect strings of lateral devices in series to achieve a TPV monolithic interconnected module (MIM) [5-7,16,17]. Fig. 3b shows a cross-sectional TEM micrograph of a LMM In_{0.69}Ga_{0.31}As TPV structure grown on an InP substrate with a metamorphic $InAs_yP_{1-y}$ buffer by MBE that is representative of our TPV growths regardless of which window layer material (InAs_{0.32}P_{0.68} or In_{0.7}Al_{0.3}As) is used. The misfit and threading dislocations are predominantly contained in the buffer layer, and plan-view TEM measurements previously revealed threading dislocation densities on the order of $2-4 \times 10^6$ cm⁻² within the active device layers [14]. The thickness and doping level in the n-type In_{0.69}Ga_{0.31}As emitter were optimized by growing several n/p $In_{0.69}Ga_{0.31}As$ diode structures on $InAs_{\nu}P_{1-\nu}/InP$ substrates. Ti/Au (200 Å/3 µm) metallization was used for both front and back ohmic contacts, and a silicon nitride (Si₃N₄) dielectric layer was deposited using plasma enhanced chemical vapor deposition to prevent the interconnect metallization from short-circuiting the individual cells in the MIM configuration. A similar Si₃N₄ layer was used on the top surface of the In_{0.7}Al_{0.3}As window layer after removal of the heavily doped In_{0.69}Ga_{0.31}As cap, to act as an antireflection coating (ARC). On the backside of the semi-insulating InP substrate, a Si₃N₄/Au coating was used as a back surface reflector to reflect the light not absorbed on the first pass and to recuperate sub-bandgap light at the radiator. A full description of the TPV device design and processing can be found in earlier papers [2,3].

3. Results and discussion

The electrical quality of the single-junction (SJ) TPV cells was evaluated using a quartz halogen lamp assembly to measure short-circuit current (I_{sc}), open-circuit voltage (V_{oc}), and fill factor (FF) to determine the maximum output power ($P_{max} = V_{oc} * I_{sc} * FF$). The light source is a quartz halogen tungsten lamp whose spectral emission fits a graybody spectrum with a temperature of 1920 K (emissivity = 0.0230) and 2050 K (emissivity = 0.0252) with the intensity at the cell being controlled (diminished) by reducing the solid angle of illumination viewed by the cell. Fig. 4a and b shows I-V characteristics obtained at 1920 K and 2050 K, respectively, from these single-junction TPV cells. Short-circuit current density $(J_{\rm sc})$, open-circuit voltage $(V_{\rm oc})$, fill factor (FF) and output power $(P_{\rm max})$ values of 1.11 A/cm², 334 mV, 65.6% and 6.9 mW, respectively, were obtained for a SJ cell at 1920 K and J_{sc} , V_{oc} , FF and P_{max} values of 2.26 A/cm², 355 mV, 66.5% and 15 mW (0.532 W/cm²), respectively, were obtained from the same cell measured at 2050 K. These V_{oc} values are comparable to those obtained from similar 0.6 eV bandgap TPV cells grown by MOCVD and MBE that employed more typical InAs_{0.32}P_{0.68} window layer materials [2,5–7].

In order to investigate the dominant current transport mechanisms for this SJ TPV cell, a series of J-V measurements were performed by varying light intensities to measure J_{sc} , V_{oc} , and FF. Fig. 5 shows J_{sc} versus V_{oc} data obtained for a SJ TPV device, where an increase in J_{sc} corresponds to an increase in light intensity for the J-V measurements. The value of V_{oc} increases logarithmically with J_{sc} , as expected, and by performing a linear fit of the data in Fig. 5 and using the ideal diode equation [1]

$$J_{\rm sc} = J_{\rm o} \bigg\{ \exp\left(\frac{qV_{\rm oc}}{nkT_{\rm cell}}\right) - 1 \bigg\},\tag{1}$$

where *n*, the diode ideality factor and J_0 , the dark saturation current density, can be estimated [4]. In this expression, k is the Boltzmann



Area = 2.82×10^{-2} cm²

1920 K (emissivity = 0.0230)

Fig. 4. Current density versus voltage characteristics of SJ In_{0.69}Ga_{0.31}As TPV cell at (a) 1920 K and (b) 2050 K, respectively. The light source is a quartz halogen tungsten lamp whose spectral emission matches a graybody spectrum with a temperature of 1920 K (emissivity = 0.0230) and 2050 K (emissivity = 0.0252).



Fig. 5. Variation of short-circuit current density (J_{sc}) with open-circuit voltage (V_{oc}) and fill factor (FF) with J_{sc} for a SJ TPV cell with $E_g = 0.60$ eV obtained as a function of incident light intensity. The fill factor value approaches ${\sim}66\%$ at higher illumination intensities

constant and T_{Cell} is the temperature of the cell during measurement (kept constant at 298 K). The data implies a change in the dominant transport mechanism from one that is likely to be depletion region recombination-limited at lower light intensities (denoted by the $n \sim 1.8$ region in the figure) to one that is more controlled by injection/diffusion in the higher injection regime, which is where TPV devices typically operate. In this higher injection regime, a linear fit yields $n \approx 1.2$, with the extracted value for J_{o} determined to be 3.3×10^{-5} A/cm². The fill factor as a function of $J_{\rm sc}$ is also shown in Fig. 5 and approaches ~66% at higher illumination intensities. These performance values were observed with a high degree of uniformity ($\pm 1\%$ of cell parameters J_{sc} , V_{oc} , FF and P_{max}) across a 2 in. InP wafer under identical illumination conditions and intensities (1920 K graybody spectrum; emissivity of 0.0230). This is particularly noteworthy for $V_{\rm oc}$ and FF, which are sensitive to the impact of small variations in material uniformity on device recombination and shunting. The high performance and uniformity of the devices implies very high quality growth of the In-AlAs window and of the InAlAs/InGaAs window/emitter interface.

Internal quantum efficiency (IQE) measurements were used to further investigate the detailed carrier collection efficiency for the In_{0.69}Ga_{0.31}As/In_{0.7}Al_{0.3}As emitter/window heterostructures and of the entire TPV cell structure. The IQE data were obtained using the measured external quantum efficiency and reflectance from the SJ TPV cell. As seen in Fig. 6, uniform and very high $(\sim 97\%)$ quantum efficiencies are obtained from a SJ In_{0.69}Ga_{0.31}As TPV cell for photon wavelengths in excess of \sim 1.4 μ m. This illustrates the effectiveness of the In_{0.69}Ga_{0.31}As n/p junction in collecting photo-generated minority carrier electron throughout the p-type base layer. It also indicates a high carrier lifetime and a diffusion length in excess of the 2.5 µm base thickness for minority carrier electrons in the p-type In_{0.69}Ga_{0.31}As base layer. However, in the shorter wavelength region, which is more sensitive to recombination within the emitter/window interface, an even higher value of IQE is maintained for this SJ TPV cell, approaching \sim 100%. These results demonstrate the general viability of using In-AlAs window lavers as an alternative for InAsP within InGaAsbased TPV devices. More specifically for MBE growth, by substituting InAlAs/InGaAs/InAsP heterostructure design for the more conventional symmetric InAsP/InGaAs/InAsP heterostructure, known growth transition issues [2] at the window/emitter interface that are problematic due to switching from a non-P containing emitter layer (InGaAs) to a P-containing window layer (InAsP), the conventional double heterostructure design have been completely avoided.



Fig. 6. Internal quantum efficiency of a SJ TPV cell with E_g = 0.60 eV. The band gap, 1.005 eV (Ref. [12]) of In_{0.7}Al_{0.3}As window layer is indicated in this figure.

4. Conclusion

This paper describes the application of $In_{0.7}AI_{0.3}As$ window layer material on a lattice-mismatched $In_{0.69}Ga_{0.31}As$ TPV cells grown by MBE. High performance devices were obtained, with SJ cells displaying a V_{oc} of 0.355 V, FF of 66.5% and power density of 0.532 W/cm² measured at a J_{sc} value of 2.26 A/cm². Internal quantum efficiency measurements revealed outstanding short wavelength carrier collection efficiency, which is attributed to the avoidance of a growth interruption at the critical window/emitter interface that is achieved by using an $In_{0.7}AI_{0.3}As/In_{0.69}Ga_{0.31}As$ heterostructure for these lattice-mismatched MBE-grown devices. This demonstrates great promise for the use of InAlAs as an alternative window layer for InGaAs-based TPV devices.

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